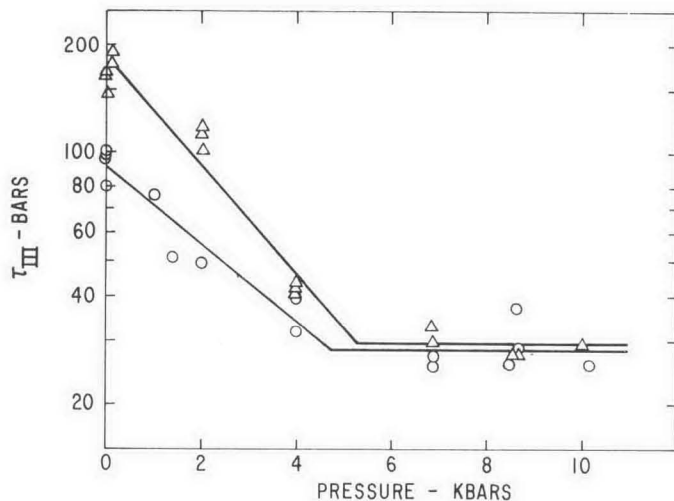
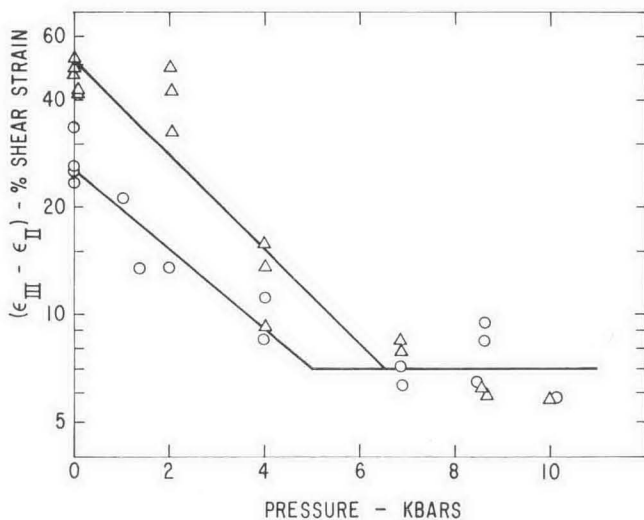


Fig. 2



Dependence of the stress for initiation of stage III,  $\tau_{III}$ , on pressure; circles refer to  $\dot{\epsilon} = 5 \times 10^{-4} \text{ sec}^{-1}$  and triangles to  $\dot{\epsilon} = 1.2 \times 10^{-2} \text{ sec}^{-1}$ . The straight lines between 1 atm and 4 kb are fitted by least squares; the slopes are given in the text.

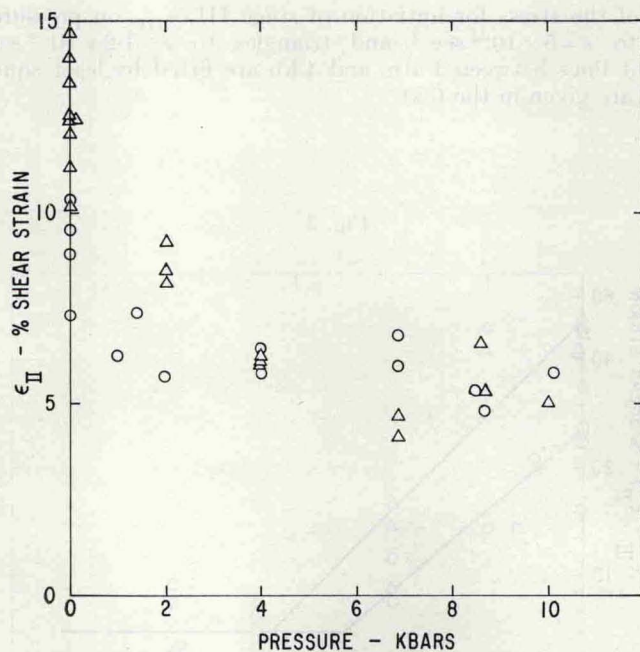
Fig. 3



Dependence of the range of stage II,  $(\epsilon_{III} - \epsilon_{II})$ , on pressure, where symbols are as for fig. 2. The sloping lines are fitted by least squares; above 6 kb a simple means of all the data points is plotted.

In figs. 2 and 3 it also is apparent that the effect on  $P$  of the initiation of stage III saturates in the vicinity of 5 kb for both strain rates. According to our qualitative association of decreasing  $\tau_{III}$  with increasing  $\gamma$ , noted in the introduction, it would follow that either the increase of  $\gamma$  saturates or the stacking-fault width becomes so narrow it is no longer sensitive to increasing  $\gamma$ . In fig. 2, which is of most interest, the horizontal lines represent the mean values (from 5 to 10 kb) of  $\tau_{III}$  at each strain rate. It is also readily apparent from fig. 2 that the strain rate sensitivity of  $\tau_{III}$  decreases greatly with increase in pressure. Extracting points from the solid lines shown one finds that  $(\partial \ln \tau_{III} / \partial \ln \dot{\epsilon}) = 0.22$  at 1 atm (in good agreement with the 1 atm data of Hesse) then decreases to 0.099 at 4 kb and finally becomes roughly constant at 0.01 between 5 and 10 kb. The last value is probably only reliable within a factor of 2 or 3. A far greater number of samples would have to be tested to fix it more accurately. It is clear, however, that a strong decrease of  $(\partial \ln \tau_{III} / \partial \ln \dot{\epsilon})$  with pressure is well established.

Fig. 4



Dependence of the range of stage I,  $\epsilon_{II}$ , on pressure; circles are for  $\dot{\epsilon} = 5 \times 10^{-4} \text{ sec}^{-1}$  and triangles for  $\dot{\epsilon} = 1.2 \times 10^{-2} \text{ sec}^{-1}$ .

On comparison of the present  $\tau$ - $\epsilon$  curves with those reported by Aladag *et al.* reasonable agreement is found. The stronger dependence of  $\tau_{III}$  and  $(\epsilon_{III} - \epsilon_{II})$  on  $P$  noted here is, in considerable measure, due to the 1%